Oscillation criteria for third-order semi-canonical differential equations with unbounded neutral coefficients

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Abstract. In this paper, we investigate the oscillatory behavior of solutions to a class of third-order differential equations of the form

$$\mathcal{L}z(t) + f(t)y^{\beta}(\sigma(t)) = 0,$$

where $\mathcal{L}z(t) = (p(t)(q(t)z'(t))')'$ is a semi-canonical operator and $z(t) = y(t) + g(t)y(\tau(t))$. The main idea is to convert the semi-canonical operator into canonical form and then obtain some new sufficient conditions for the oscillation of all solutions. The obtained results essentially improve and complement to the known results. Examples are provided to illustrate the main results.

Mathematics Subject Classification (2010): 34C10, 34K11, 34K40.

Keywords: Oscillation, third-order, semi-canonical, unbounded neutral coefficients.

1. Introduction

In this paper, we are concerned with the oscillation of solutions of the semicanonical third-order neutral differential equation

$$\mathcal{L}z(t) + f(t)y^{\beta}(\sigma(t)) = 0, \quad t \ge t_0 > 0, \tag{1.1}$$

where \mathcal{L} is the differential operator defined by

$$\mathcal{L}z(t) = (p(t)(q(t)z'(t))')', \quad z(t) = y(t) + g(t)y(\tau(t)),$$

Received 19 September 2021; Accepted 20 January 2022.

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and β is the ratio of odd positive integers. Throughout the paper, and without further mention, we will always assume that:

- (H_1) $f, g \in C([t_0, \infty), \mathbb{R}), g(t) \ge 1, g(t) \not\equiv 1$ for large t, and $f(t) \ge 0$ is not identically zero for large t,
- (H_2) $\tau, \sigma \in C^1([t_0, \infty), \mathbb{R}), \tau(t) \leq t, \tau$ is strictly increasing, σ is nondecreasing, and $\lim_{t\to\infty} \tau(t) = \lim_{t\to\infty} \sigma(t) = \infty;$
- (H_3) the operator \mathcal{L} is in semi-canonical form, that is,

$$\int_{t_0}^{\infty} \frac{1}{p(t)} dt < \infty \text{ and } \int_{t_0}^{\infty} \frac{1}{q(t)} dt = \infty,$$

where $p, q \in C([t_0, \infty), (0, \infty))$.

By a solution of (1.1), we mean a function $y \in C([t_y, \infty), \mathbb{R})$ for some $t_y \geq t_0$ such that $z \in C^1([t_y, \infty), \mathbb{R})$, $qz' \in C^1([t_y, \infty), \mathbb{R})$, $p(qz')' \in C^1([t_y, \infty), \mathbb{R})$ and y satisfies (1.1) on $[t_y, \infty)$. We only consider those solutions of (1.1) that exist on some half-line $[t_y, \infty)$ and satisfy the condition

$$\sup\{|y(t)|: T_1 \le t < \infty\} > 0 \text{ for any } T_1 \ge t_y;$$

we tacitly assume that (1.1) possesses such solutions. Such a solution y(t) of (1.1) is said to be oscillatory if it has arbitrarily large zeros on $[t_y, \infty)$, and it is called nonoscillatory otherwise. Equation (1.1) is called oscillatory if all its solutions are oscillatory.

In the recent years many papers appeared in the literature dealing with the oscillatory and asymptotic behavior of solutions of various classes of third-order neutral type differential equations; see for example [1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 12, 13, 16, 17, 19, 20] and the references cited therein. However, except for the papers [5, 6, 12, 19, 20], all the papers mentioned above were dealing with the case when g(t) is bounded, that is, the cases when $0 \le g(t) \le g_0 < 1$, $-1 < g_0 \le g(t) \le 0$ and $0 < g(t) \le g_0 < \infty$ were studied and so the criteria obtained in these papers cannot be applied to the case $g(t) \to \infty$ as $t \to \infty$.

Moreover, very recently in [5, 6, 20] the authors studied equation (1.1) and obtained oscillation criteria where $q(t) \equiv 1$ and $p(t) \equiv 1$ or $\int_{t_0}^{\infty} \frac{1}{p(t)} dt = \infty$. Based on these observations, the aim of this paper is to obtain some oscillation criteria that can be applied not only to the case where $g(t) \to \infty$ as $t \to \infty$ but also to the cases when g(t) is bounded, $\int_{t_0}^{\infty} \frac{1}{p(t)} dt < \infty$ and $\int_{t_0}^{\infty} \frac{1}{q(t)} dt = \infty$. The main idea is to connect the semi-canonical equation (1.1) with that of canonical equations and then we obtain oscillation criteria for (1.1).

In the sequel, we deal only with positive solutions of (1.1), since if y(t) is a solution of (1.1), then -y(t) is also a solution.

2. Main results

Throughout the paper we employ the following notations:

$$A(t) := \int_t^\infty \frac{1}{p(s)} ds, \quad a(t) := p(t) A^2(t), \quad b(t) := \frac{q(t)}{A(t)},$$

Oscillation criteria for third-order differential equations

$$\begin{split} F(t) &:= A(t)f(t), \quad \Pi(t) := \int_{t_0}^t \frac{1}{a(s)} ds, \quad B(t) := \int_{t_0}^t \frac{\Pi(s)}{b(s)} ds, \\ c(t) &:= \exp\left(\int_{t_1}^t \frac{\Pi(s)}{b(s)B(s)} ds\right) \quad \text{for } t \ge t_1 \quad \text{for some } t_1 \ge t_0, \\ h(t) &:= \tau^{-1}(\sigma(t)), \quad \lambda(t) := \tau^{-1}(\eta(t)), \quad \eta \in C^1([t_0, \infty), \mathbb{R}), \\ \psi_1(t) &:= \frac{1}{g(\tau^{-1}(t))} \left[1 - \frac{c(\tau^{-1}(\tau^{-1}(t)))}{g(\tau^{-1}(\tau^{-1}(t)))c(\tau^{-1}(t))}\right], \\ \psi_2(t) &:= \frac{1}{g(\tau^{-1}(t))} \left[1 - \frac{1}{g(\tau^{-1}(\tau^{-1}(t)))}\right], \end{split}$$

and

$$R(t) := \int_{h(t)}^{\lambda(t)} \left(\frac{1}{b(u)} \int_{u}^{\lambda(t)} \frac{1}{a(v)} dv\right) du.$$

In order to ensure the nonnegativity of $\psi_1(t)$, we assume the following condition also holds:

 (H_4) There exists a $t_1 \in [t_0, \infty)$ such that

$$\frac{c(\tau^{-1}(\tau^{-1}(t)))}{g(\tau^{-1}(\tau^{-1}(t)))c(\tau^{-1}(t))} \le 1 \quad \text{for all } t \ge t_1.$$
(2.1)

Theorem 2.1. Assume that

$$\int_{t_0}^{\infty} \frac{1}{b(t)} dt = \infty.$$
(2.2)

Then the semi-canonical operator $\mathcal L$ has the following unique canonical representation

$$\mathcal{L}z(t) = \frac{1}{A(t)} \left(p(t)A^2(t) \left(\frac{q(t)}{A(t)} z'(t) \right)' \right)'.$$
(2.3)

Proof. Direct calculation shows that

$$\left(p(t)A^{2}(t)\left(\frac{q(t)}{A(t)}z'(t)\right)'\right)' = (A(t)p(t)(q(t)z'(t))' + q(t)z'(t))'$$
$$= A(t)(p(t)(q(t)z'(t))')'.$$

Therefore

$$\frac{1}{A(t)} \left(p(t)A^2(t) \left(\frac{q(t)}{A(t)} z'(t) \right)' \right)' = (p(t)(q(t)z'(t))')'.$$

Taking (2.2) into account, we see that

$$\int_{t_0}^{\infty} \frac{A(t)}{q(t)} dt = \infty,$$

and since

$$\int_{t_0}^{\infty} \frac{1}{p(t)A^2(t)} dt = \lim_{t \to \infty} \left(\frac{1}{A(t)} - \frac{1}{A(t_0)} \right) = \infty,$$

we say that (2.3) is in the canonical form. However, Trench proved in [18] that there exists only one canonical representation of \mathcal{L} (up to multiplicative constants with product 1) and so our canonical form is unique. This completes the proof.

From Theorem 2.1, it follows that (1.1) can be written in the canonical form as

$$(a(t)(b(t)z'(t))')' + F(t)y^{\beta}(\sigma(t)) = 0$$
(2.4)

and the next result is immediate.

Theorem 2.2. Assume that (2.2) holds. Then semi-canonical equation (1.1) possesses solution y(t) if and only if canonical equation (2.4) has the solution y(t).

Corollary 2.3. Assume that (2.2) holds. Then semi-canonical differential equation (1.1) has an eventually positive solution if and only if canonical equation (2.4) has an eventually positive solution.

Corollary 2.3 clearly simplifies investigation of (1.1) since for (2.4) if y(t) is an eventually positive solution, then the corresponding function z(t) satisfies either

- (I) z(t) > 0, b(t)z'(t) > 0, a(t)(b(t)z'(t))' > 0, (a(t)(b(t)z'(t))')' < 0, or
- (II) z(t) > 0, b(t)z'(t) < 0, a(t)(b(t)z'(t))' > 0, (a(t)(b(t)z'(t))')' < 0 for sufficiently large t.

Lemma 2.4. Assume that z(t) satisfies case (I) for all $t \ge t_1$ for some $t_1 \ge t_0$. Then

$$z'(t) \ge \frac{\Pi(t)}{b(t)} a(t) (b(t) z'(t))',$$
(2.5)

$$z(t) \ge B(t)a(t)(b(t)z'(t))',$$
 (2.6)

$$z(t) \ge \frac{B(t)}{\Pi(t)} b(t) z'(t), \qquad (2.7)$$

and

$$\frac{z(t)}{c(t)}$$
 is nonincreasing (2.8)

for all $t \geq t_1$.

Proof. Since a(t)(b(t)z'(t))' is positive and decreasing, we see that

$$b(t)z'(t) = b(t_1)z'(t_1) + \int_{t_1}^t a(s)\frac{(b(s)z'(s))'}{a(s)}ds$$

or

$$z'(t) \ge \frac{a(t)}{b(t)} (b(t)z'(t))' \Pi(t),$$

i.e., (2.5) holds. Integrating the last inequality from t_1 to t yields

$$z(t) \ge a(t)(b(t)z'(t))' \int_{t_1}^t \frac{\Pi(s)}{b(s)} ds = B(t)a(t)(b(t)z'(t))',$$

i.e., (2.6) holds. From (2.5), we see that $b(t)z'(t)/\Pi(t)$ is decreasing for $t \ge t_2$ for some $t_2 \ge t_1$, and therefore

$$z(t) = z(t_2) + \int_{t_2}^t \frac{b(s)z'(s)}{\Pi(s)} \frac{\Pi(s)}{b(s)} ds \ge \frac{B(t)}{\Pi(t)} b(t)z'(t).$$

From the last inequality, we see that

$$\left(\frac{z(t)}{c(t)}\right)' = \frac{\left(z'(t) - \frac{\Pi(t)}{b(t)B(t)}z(t)\right)}{c(t)} \le 0$$

for $t \ge t_3$ for some $t_3 \ge t_2$. Hence, z(t)/c(t) is non-increasing. This completes the proof.

Theorem 2.5. Let (2.2) holds. Assume that there exists a nondecreasing function $\eta \in C^1([t_0,\infty),\mathbb{R})$ such that $\sigma(t) \leq \eta(t) < \tau(t)$ for all $t \geq t_0$. If both first-order delay differential equations

$$X'(t) + F(t)\Psi_1^{\beta}(\sigma(t))B^{\beta}(h(t))X^{\beta}(h(t)) = 0$$
(2.9)

and

$$W'(t) + F(t)\Psi_{2}^{\beta}(\sigma(t))R^{\beta}(t)W^{\beta}(\lambda(t)) = 0$$
(2.10)

oscillate, then (1.1) oscillates.

Proof. Let y(t) be a nonoscillatory solution of equation (1.1), say y(t) > 0, $y(\tau(t)) > 0$, and $y(\sigma(t)) > 0$ for $t \ge t_1$ for some $t_1 \ge t_0$. From Corollary 2.3, y(t) is also a positive solution of (2.4) for $t \ge t_1$. Then the corresponding function z(t) satisfies either case (I) or case (II) for $t \ge t_2$ for some $t_2 \ge t_1$.

First, we consider case (I). From the definition of z, we get

$$y(t) = \frac{1}{g(\tau^{-1}(t))} \left[z(\tau^{-1}(t)) - y(\tau^{-1}(t)) \right]$$

$$\geq \frac{z(\tau^{-1}(t))}{g(\tau^{-1}(t))} - \frac{z(\tau^{-1}(\tau^{-1}(t)))}{g(\tau^{-1}(\tau^{-1}(t)))}.$$
(2.11)

Now $\tau(t) \leq t$ and τ is strictly increasing, so τ^{-1} is increasing and $t \leq \tau^{-1}(t)$. Thus,

$$\tau^{-1}(t) \le \tau^{-1}(\tau^{-1}(t))$$

From this and the fact that z(t)/c(t) is nonincreasing, we see that

$$z\left(\tau^{-1}(\tau^{-1}(t))\right) \le \frac{c(\tau^{-1}(\tau^{-1}(t)))z(\tau^{-1}(t))}{c(\tau^{-1}(t))}.$$
(2.12)

Using (2.12) in (2.11) yields

$$y(t) \ge \psi_1(t)z(\tau^{-1}(t)).$$
 (2.13)

Since $\lim_{t\to\infty} \sigma(t) = \infty$, we can choose $t_3 \ge t_2$ such that $\sigma(t) \ge t_2$ for all $t \ge t_3$. Thus, it follows from (2.13) that

$$y(\sigma(t)) \ge \psi_1(\sigma(t))z(h(t)) \quad \text{for } t \ge t_3.$$
(2.14)

Combining (2.14) with (2.4) yields

$$(a(t)(b(t)z'(t))')' + F(t)\psi_1^\beta(\sigma(t))z^\beta(h(t)) \le 0 \quad \text{for } t \ge t_3.$$
(2.15)

From (2.6), we have

$$z(h(t)) \ge B(h(t))a(h(t))(b(h(t))z'(h(t)))'.$$
(2.16)

Using (2.16) in (2.15) and letting X(t) = a(t)(b(t)z'(t))', we see that X(t) is a positive solution of the first-order delay differential inequality

$$X'(t) + F(t)\psi_1^{\beta}(\sigma(t))B^{\beta}(h(t))X^{\beta}(h(t)) \le 0.$$
(2.17)

Therefore, by Corollary 1 of [14], we conclude that (2.9) also has a positive solution, which is a contradiction.

Next, we consider case (II). Since z is strictly decreasing and $\tau(t) \leq t$, we have

$$z(\tau^{-1}(t)) \ge z(\tau^{-1}(\tau^{-1}(t)))$$

and using this in (2.11), we obtain

$$y(t) \ge \psi_2(t) z(\tau^{-1}(t)).$$

Hence,

$$y(\sigma(t)) \ge \psi_2(\sigma(t))z(h(t)) \tag{2.18}$$

for $t \ge t_3$ for some $t_3 \ge t_2$. Using (2.18) in (2.4) yields

$$(a(t)(b(t)z'(t))')' + F(t)\psi_2^\beta(\sigma(t))z^\beta(h(t)) \le 0 \quad \text{for } t \ge t_3.$$
(2.19)

For $t \geq s \geq t_3$, we have

$$b(t)z'(t) - b(s)z'(s) = \int_{s}^{t} \frac{a(u)(b(u)z'(u))'}{a(u)} du,$$

or

$$-z'(s) \ge \left(\frac{1}{b(s)} \int_s^t \frac{1}{a(u)} du\right) a(t)(b(t)z'(t))'$$

Again integrating, we have

$$-z(t) + z(s) \ge \left(\int_s^t \frac{1}{b(u)} \left(\int_u^t \frac{1}{a(v)} dv\right) du\right) a(t)(b(t)z'(t))',$$

or

$$z(s) \ge \left[\int_{s}^{t} \frac{1}{b(u)} \left(\int_{u}^{t} \frac{1}{a(v)} dv\right) du\right] a(t)(b(t)z'(t))'.$$
(2.20)

Since $\sigma(t) \leq \eta(t)$ and the fact that τ is strictly increasing, we have

$$\tau^{-1}(\sigma(t)) \le \tau^{-1}(\eta(t)).$$

Setting $s = \tau^{-1}(\sigma(t))$ and $t = \tau^{-1}(\eta(t))$ into (2.20), we obtain

$$z(h(t)) \ge \left(\int_{h(t)}^{\lambda(t)} \frac{1}{b(u)} \left(\int_{u}^{\lambda(t)} \frac{1}{a(v)} dv\right) du\right) a(\lambda(t))(b(\lambda(t))z'(\lambda(t)))'.$$
(2.21)

Using (2.21) in (2.19) and letting W(t) = a(t)(b(t)z'(t))', we see that W is a positive solution of the first-order delay differential inequality

$$W'(t) + F(t)\psi_2^\beta(\sigma(t))R^\beta(t)W^\beta(\lambda(t)) \le 0.$$
(2.22)

The remaining part of the proof is similar to the case (I) and hence the details are not repeated. This completes the proof. $\hfill \Box$

Corollary 2.6. Let (2.2) holds and $\beta = 1$. Assume that there exists a nondecreasing function $\eta \in C^1([t_0, \infty), \mathbb{R})$ such that $\sigma(t) \leq \eta(t) < \tau(t)$ for all $t \geq t_0$. If

$$\liminf_{t \to \infty} \int_{h(t)}^{t} F(s)\psi_1(\sigma(s))B(h(s))ds > \frac{1}{e}$$
(2.23)

and

$$\liminf_{t \to \infty} \int_{\lambda(t)}^{t} F(s)\psi_2(\sigma(s))R(s)ds > \frac{1}{e},$$
(2.24)

then (1.1) is oscillatory.

Proof. The proof follows from a well-known result in [11] and Theorem 2.5, and hence the details are omitted. \Box

Corollary 2.7. Let (2.2) holds and $0 < \beta < 1$. Assume that there exists a nondecreasing function $\eta \in C^1([t_0, \infty), \mathbb{R})$ such that $\sigma(t) \leq \eta(t) < \tau(t)$ for all $t \geq t_0$. If

$$\int_{T}^{\infty} F(t)\psi_{1}^{\beta}(\sigma(t))B^{\beta}(h(t))dt = \infty$$
(2.25)

and

or

$$\int_{T}^{\infty} F(t)\psi_{2}^{\beta}(\sigma(t))R^{\beta}(t)dt = \infty$$
(2.26)

for all sufficiently large $T \in [t_0, \infty)$ with $\sigma(t) \ge t_0$ for all $t \ge T$, then (1.1) oscillates.

Proof. Proceeding exactly as in the proof of Theorem 2.5, we again arrive at (2.17) and (2.22) for $t \ge t_3$. Since h(t) < t and X(t) is positive and decreasing, inequality (2.17) takes the form

$$X'(t) + F(t)\psi_{1}^{\beta}(\sigma(t))B^{\beta}(h(t))X^{\beta}(t) \leq 0,$$

$$\frac{X'(t)}{X^{\beta}(t)} + F(t)\psi_{1}^{\beta}(\sigma(t))B^{\beta}(h(t)) \leq 0.$$
 (2.27)

Integrating (2.27) from t_3 to t yields

$$\int_{t_3}^t F(s) \psi_1^\beta(\sigma(s)) B^\beta(h(s)) ds \leq \frac{X^{1-\beta}(t_3)}{1-\beta} < \infty \ \text{ as } t \to \infty$$

which contradicts (2.25). The remainder of the proof follows from $\lambda(t) < t$ and inequality (2.22). The proof is complete.

In our final result, assume that $\sigma(t) = t - \delta_1$, $\tau(t) = t - \delta_3$ and $\eta(t) = t - \delta_2$, where δ_1 , δ_2 and δ_3 are positive real numbers.

Corollary 2.8. Let (2.2) holds and $\beta > 1$. If $\delta_1 \ge \delta_2 > \delta_3$,

$$\liminf_{t \to \infty} \beta^{-t/(\delta_1 - \delta_3)} \log \left(F(t) \psi_1^\beta(t - \delta_1) B^\beta(t + \delta_3 - \delta_1) \right) > 0$$
(2.28)

and

$$\liminf_{t \to \infty} \beta^{-t/(\delta_2 - \delta_3)} \log \left(F(t) \psi_2^\beta(t - \delta_1) R^\beta(t) \right) > 0, \tag{2.29}$$

then (1.1) oscillates.

Proof. Application of (2.28) and (2.29) and Corollary 1.2 of [15] imply that (2.9) and (2.10) oscillate. Hence, by Theorem 2.5, equation (1.1) oscillates.

3. Examples

In this section, we present some examples to show the importance of the main results.

Example 3.1. Consider the third-order linear neutral differential equation

$$\left(t^2 \left(\frac{1}{t} \left(y(t) + 16y\left(\frac{t}{2}\right)\right)'\right)' + \frac{f_0}{t^2} y\left(\frac{t}{4}\right) = 0, \quad t \ge 1.$$

$$(3.1)$$

Here $p(t) = t^2$, q(t) = 1/t, g(t) = 16, $f(t) = f_0/t^2$ with $f_0 > 0$, $\tau(t) = t/2$, $\sigma(t) = t/4$ and $\beta = 1$. Then A(t) = 1/t, a(t) = 1, b(t) = 1, $F(t) = f_0/t^3$ and the transformed equation is

$$\left(y(t) + 16y\left(\frac{t}{2}\right)\right)^{\prime\prime\prime} + \frac{f_0}{t^3}y\left(\frac{t}{4}\right) = 0, \quad t \ge 1,$$
(3.2)

which is in canonical form. Simple calculation show that

$$\Pi(t) = t - 1$$
, $B(t) = (t - 1)^2/2$, $c(t) = (t - 1)^2$, and $\psi_2(t) = 15/256$.

Since (2.1) holds, we have $\psi_1(t) \ge 0$ and

$$\psi_1(t) = \frac{1}{16} \left[1 - \frac{(4t-1)^2}{16(2t-1)^2} \right] \ge \frac{7}{256}$$

By choosing $\eta(t) = t/3$, we see that h(t) = t/2, $\lambda(t) = 2t/3$ and $R(t) = t^2/72$. It is clear that condition (2.2) holds. Condition (2.23) becomes

$$\liminf_{t \to \infty} \int_{t/2}^t \frac{f_0}{2^9} \left(\frac{3}{s} - \frac{14}{s^2} + \frac{15}{s^3}\right) ds = \frac{3f_0 \ln 2}{2^9},$$

and so condition (2.22) is satisfied if $f_0 > \frac{2^9}{3e \ln 2}$.

Condition (2.24) becomes

$$\liminf_{t \to \infty} \int_{2t/3}^t \frac{5f_0}{3 \times 2^{11}} \frac{1}{s} ds = \frac{5f_0 \ln 3/2}{3 \times 2^{11}} \frac{1}{s} ds$$

that is, (2.24) is satisfied if $f_0 > \frac{3 \times 2^{11}}{5e \ln 3/2}$. Thus, by Corollary 2.6, equation (3.1) is oscillatory if $f_0 > \frac{3 \times 2^{11}}{5e \ln 3/2}$.

Note that canonical equation (3.2) is considered in [20] and proved that (3.2) is oscillatory if $f_0 > \frac{3 \times 2^{11}}{5 \ln 3/2}$. Hence, Corollary 2.6 improves Theorem 2.7 of [20].

Example 3.2. Consider the third-order sublinear neutral differential equation

$$\left(t^2 \left(\frac{1}{t} \left(y(t) + ty\left(\frac{t}{2}\right)\right)'\right)' + \frac{f_0}{t^\alpha} y^{3/5} \left(\frac{t}{10}\right) = 0, \quad t \ge 16.$$

$$(3.3)$$

Here $p(t) = t^2$, q(t) = 1/t, g(t) = t, $f(t) = f_0/t^{\alpha}$ with $f_0 > 0$ and $\alpha \le 3/5$, $\tau(t) = t/2$, $\sigma(t) = t/10$ and $\beta = 3/5$. Then A(t) = 1/t, a(t) = 1, b(t) = 1, $F(t) = f_0/t^{\alpha+1}$ and the transformed equation is

$$\left(y(t) + ty\left(\frac{t}{2}\right)\right)^{\prime\prime\prime} + \frac{f_0}{t^{\alpha+1}}y^{3/5}\left(\frac{t}{10}\right) = 0, \quad t \ge 16, \tag{3.4}$$

which is in canonical form. Simple calculation shows that

$$\Pi(t) = t - 16, \ B(t) = (t - 16)^2/2, \ c(t) = (t - 16)^2, \ \text{and} \ \psi_2(t) = \frac{4t - 1}{8t^2} > 0.$$

Since (2.1) holds, we have $\psi_1(t) \ge 0$ and $\psi_1(t) \ge \frac{4t-9}{8t^2}$. By choosing $\eta(t) = t/8$, we see that h(t) = t/5, $\lambda(t) = t/4$ and $R(t) = t^2/800$. It is clear that condition (2.2) holds. For any $T \ge t_0$ with $\sigma(t) \ge t_0$, condition (2.25) becomes

$$\int_{T}^{\infty} \frac{f_0}{t^{\alpha+1}} \left(\frac{10t-225}{2t^2}\right)^{3/5} \left(\frac{t-80}{\sqrt{50}}\right)^{6/5} dt \ge d_1 \int_{T_1}^{\infty} \frac{1}{t^{\alpha+2/5}} dt = \infty,$$

where $d_1 > 0$ is a constant and $T_1 \ge T$.

Condition (2.26) becomes

$$\int_{T}^{\infty} \frac{f_0}{t^{\alpha+1}} \left(\frac{10t-25}{2t^2}\right)^{3/5} \left(\frac{t^2}{800}\right)^{3/5} dt \ge d_2 \int_{T_1}^{\infty} \frac{1}{t^{\alpha+2/5}} = \infty,$$

where $d_2 > 0$ is a constant and $T_1 \ge T$. Thus, by Corollary 2.7, equation (3.3) is oscillatory if $\alpha \le 3/5$.

Note that canonical equation (3.4) is considered in [20] and proved that (3.4) is oscillatory if $\alpha = \frac{1}{5}$. Hence, Corollary 2.7 improves Theorem 2.8 of [20].

Example 3.3. Consider the third-order superlinear neutral differential equation

$$\left(t^2 \left(\frac{1}{t} \left(y(t) + ty \left(t - 2\right)\right)'\right)' + t \exp(4^t) y^3(t - 4) = 0, \quad t \ge 2.$$
(3.5)

Here $p(t) = t^2$, q(t) = 1/t, g(t) = t, $f(t) = t \exp(4^t)$, $\tau(t) = t - 2$, $\sigma(t) = t - 4$ and $\beta = 3$. Then A(t) = 1/t, a(t) = 1, b(t) = 1, $F(t) = \exp(4^t)$ and the transformed equation is

$$(y(t) + ty(t-2))''' + \exp(4^t)y^3(t-4) = 0,$$
(3.6)

which is in canonical form. A simple calculation show that

$$\Pi(t) = t - 2, \ B(t) = (t - 2)^2/2, \ c(t) = (t - 2)^2/2$$

$$\psi_1(t) = \frac{1}{t+2} \left[1 - \frac{(t+2)^2}{(t+4)t^2} \right] \ge \frac{t}{(t+2)(t+4)} \ge 0 \text{ and } \psi_2(t) = \frac{t+3}{(t+2)(t+4)} \ge 0.$$

By choosing $\eta(t) = t - 3$, we see that h(t) = t - 2, $\lambda(t) = t - 1$, R(t) = 1/2, $\delta_1 = 4$, $\delta_2 = 3$, $\delta_3 = 2$. As in Examples 3.1 and 3.2, it is easy to see that conditions (2.2), (2.28) and (2.29) are satisfied. Thus, by Corollary 2.8, equation (3.5) is oscillatory.

4. Conclusion

In this paper, we have established some new oscillation criteria for (1.1). The results are obtained by converting (1.1) into canonical type equation. Hence, the results are new and complement to those in [5, 6, 12, 20]. Also we have shown that the results obtained here improve those in [20].

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