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Finite time blow-up for quasilinear wave equations with nonlinear dissipation

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Abstract. In this paper we consider a class of quasilinear wave equations

$$u_{tt} - \Delta_{\alpha}u - \omega_1 \Delta u_t - \omega_2 \Delta_{\beta}u_t + \mu |u_t|^{m-2}u_t = |u|^{p-2}u,$$

associated with initial and Dirichlet boundary conditions. Under certain conditions on α, β, m, p , we show that any solution with positive initial energy, blows up in finite time. Furthermore, a lower bound for the blow-up time will be given.

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1. Introduction

In this paper, we would like to study the blow-up of solutions of the following initial boundary value problem of a quasilinear wave equation

$$\begin{cases} u_{tt} - \Delta_{\alpha} u - \omega_1 \Delta u_t - \omega_2 \Delta_{\beta} u_t + \mu |u_t|^{m-2} u_t = |u|^{p-2} u, & x \in \Omega, \quad t > 0, \\ u(x,t) = 0, & x \in \partial\Omega, \quad t > 0, \\ u(x,0) = u_0(x), & u_t(x,0) = u_1(x), & x \in \Omega. \end{cases}$$
(1.1)

Here, Ω is a bounded domain of \mathbb{R}^n with a smooth boundary $\partial\Omega$. Additionally, we assume that

$$u_0 \in W_0^{1,\alpha}(\Omega), \quad u_1 \in L^2(\Omega), \tag{1.2}$$

and $\alpha, \beta, \omega_1, \omega_2, \mu, m, p$ are positive constants, with

$$\begin{cases}
2 \alpha, \\
2
(1.3)$$

The operator Δ_{α} is the classical α -Laplacian given by:

$$\Delta_{\alpha} u = \operatorname{div} \left(|\nabla u|^{\alpha - 2} \nabla u \right)$$

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Notice that $\Delta_{\beta} u_t$ is a quasilinear strong damping term, and it is degenerate when $\beta > 2$.

Nonlinear hyperbolic equations of the type (1.1) have been investigated in the papers [2, 5, 7, 9, 15], and the references therein. Several examples of this type arise in physics, for example, the problem (1.1) represents a longitudinal motion of a viscoelastic rod obeying the nonlinear Voight model.

Zhijiang [14] proved a blow up result for the problem (1.1) when the initial energy is sufficiently negative. This result was extended by Messaoudi and Houari [8] to a situation when the solution has negative initial energy. Liu and Wang [6] studied a more general model including (1.1), and by improving the arguments in [14] and [8] they established a blow-up result in the subcritical initial energy case, i.e. E(0) < d, where E(0) is the initial energy and d is the depth of the potential well.

For $\alpha = \beta = m = 2$, equation in (1.1) reduces to the linearly damped wave equation

$$u_{tt} - \Delta u + \omega \Delta u_t + \mu u_t = |u|^{p-2} u. \tag{1.4}$$

Gazzola and Squassina [3] studied (1.4) and gave a necessary and sufficient conditions for blow-up if E(0) < d. Recently, Yang and Xu [13] gave a sufficient condition for blow-up if E(0) > d. Sun et al. [12] obtained, for (1.4), an estimate of the lower bound for the blow-up time when 2 . This work was extended by Guo and Liu $[4] to the case when the exponent <math>p \in \left(\frac{2(n-1)}{n-2}, \frac{2(n^2-2)}{n-2}\right]$. Later, in the case of $\omega > 0$, Baghaei [1] improved the results in [12] and [4] by enlarging the upper bound for pto 2^* .

In related work, Song and Xue [11] studied the following nonlinear wave equation with strong damping

$$u_{tt} - \Delta u + \int_0^t g((t-\tau)\Delta u(\tau)d\tau - \Delta u_t = |u|^{p-2}u.$$
 (1.5)

They introduced a new technique to obtain a finite time blow-up result with arbitrary high initial energy in the case of linear strong damping. By applying the technique similar to that in [11], Song [10] extended the result in [11] to the case of nonlinear weak damping $\mu |u_t|^{m-2}u_t$ in place of $-\Delta u_t$ in (1.5).

In this paper, by using the technique in [10], we give sufficient conditions for finite time blow-up of solutions of (1.1), in the case $E(0) \ge d$. Furthermore, by using the techniques in [4], we obtain a lower bound for the blow-up time.

2. Preliminaries

We denote by $\|.\|_p$ the $L^p(\Omega)$ norm $(2 \leq p < \infty)$, and by (.,.) the L^2 inner product. We introduce the following functional space

$$\begin{split} \mathcal{H} &:= L^{\infty}([0,T), W^{1,\alpha}_0(\Omega)) \cap W^{1,\infty}([0,T), L^2(\Omega)) \\ & \cap W^{1,\beta}([0,T), W^{1,\beta}(\Omega)) \cap W^{1,m}([0,T), L^m(\Omega)), \end{split}$$

for T > 0, and the energy functional

$$E(t) := \frac{1}{2} \|\nabla u\|_{\alpha}^{\alpha} + \frac{1}{2} \|u_t\|_2^2 - \frac{1}{p} \|u\|_p^p$$

We define also the following constant

$$\lambda = B_*^{-\frac{p}{p-\alpha}},$$

where B_* is the best constant of the Sobolev embedding $W_0^{1,\alpha}(\Omega) \hookrightarrow L^p(\Omega)$. Finally, we characterize the depth of the potential well d as follows:

$$d = \left(\frac{1}{\alpha} - \frac{1}{p}\right)\lambda^2.$$

Lemma 2.1. Let u be a global solution to problem (1.1). Then we have

$$E'(t) = -\omega_1 \|\nabla u_t\|_2^2 - \omega_2 \|\nabla u_t\|_{\beta}^{\beta} - \mu \|u_t\|_m^m, \quad \forall t \ge 0.$$

As a consequence, we have the following inequalities:

$$E(t) \le E(0), \quad \forall t \ge 0, \tag{2.1}$$

and

$$-E'(t) \ge \omega_1 \|\nabla u_t\|_2^2, \quad -E'(t) \ge \omega_2 \|\nabla u_t\|_{\beta}^{\beta}, \quad -E'(t) \ge \mu \|u_t\|_m^m.$$
(2.2)

Subsequently, we state the following theorems (see [6]).

Theorem 2.2 (Local existence). Assume that conditions (1.2) and (1.3) hold. Then problem (1.1) has a unique local solution $u \in \mathcal{H}$.

Theorem 2.3 (Blow-up for E(0) < d). Assume (1.2) and (1.3) hold. Assume further that $\alpha, \beta, m \ge 2$ and $p > \alpha > \max\{m, \beta\}$. Suppose E(0) < d and

$$\|\nabla u_0\|_{\alpha} > \lambda. \tag{2.3}$$

Then u blows up in finite time.

3. Finite time blow-up

In this section we extend the blow-up result in [8] to the case $E(0) \ge d$. Here is our main result:

Theorem 3.1 (Blow-up for $E(0) \ge d$). Assume (1.2), (2.3) and (1.3) hold. Assume further that $\alpha, \beta, m > 2, \alpha > \beta$ and $p > \max\{m, \alpha\}$. Suppose $E(0) \ge d$ and

$$(u_t(0), u(0)) > ME(0), \tag{3.1}$$

where M > 0 is defined in (3.7), then the solution $u \in \mathcal{H}$ of (1.1) blows up in finite time.

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Proof. Assume by contradiction that u(t) is a global solution of (1.1). Setting

$$F(t) := \frac{1}{2} \|u(t)\|_2^2,$$

it follows from (1.1) that

$$F''(t) = \|u_t\|_2^2 + \|u\|_p^p - \|\nabla u\|_{\alpha}^{\alpha} - \omega_1(\nabla u_t, \nabla u) - \omega_2(|\nabla u_t|^{\beta-2}\nabla u_t, u) - \mu(|u_t|^{m-2}u_t, u).$$
(3.2)

By using Hölder's inequality and Young's inequality, we estimate the two last terms in the right-hand side of the previous equation, as follows

$$\begin{split} (\nabla u_t, \nabla u) &\leq \eta \|\nabla u\|_2^2 + \frac{1}{4\eta} \|\nabla u_t\|_2^2, \quad \eta > 0, \\ (|\nabla u_t|^{\beta - 2} \nabla u_t, u) &\leq \frac{1}{\beta} \sigma^{\beta} \|\nabla u\|_{\beta}^{\beta} + \frac{\beta - 1}{\beta} \sigma^{\beta/(1 - \beta)} \|\nabla u_t\|_{\beta}^{\beta}, \quad \sigma > 0, \\ (|u_t|^{m - 2} u_t, u) &\leq \frac{1}{m} \delta^m \|u\|_m^m + \frac{m - 1}{m} \delta^{m/(1 - m)} \|u_t\|_m^m, \quad \delta > 0. \end{split}$$

So, that to the convexity of the function y^x/x for $y \ge 0$ and x > 0, we have

$$\frac{\delta^m}{m} \|u\|_m^m \le \frac{s}{2} \delta^m \|u\|_2^2 + \frac{1-s}{p} \delta^m \|u\|_p^p, \qquad s = \frac{p-m}{p-2},$$
$$\frac{1}{\beta} \sigma^\beta \|\nabla u\|_\beta^\beta \le \frac{\theta}{2} \sigma^\beta \|\nabla u\|_2^2 + \frac{1-\theta}{\alpha} \sigma^\beta \|\nabla u\|_\alpha^\alpha, \qquad \theta = \frac{\alpha-\beta}{\alpha-2}$$

Hence, (3.2) becomes

$$F''(t) \geq \|u_t\|_2^2 - \left[1 + \frac{\omega_2(1-\theta)}{\alpha}\sigma^{\beta}\right] \|\nabla u\|_{\alpha}^{\alpha} - \frac{\mu s}{2}\delta^m \|u\|_2^2 - \left(\omega_1\eta + \frac{\omega_2\theta}{2}\sigma^{\beta}\right) \|\nabla u\|_2^2 + \left[1 - \frac{\mu(1-s)}{p}\delta^m\right] \|u\|_p^p - \frac{\omega_1}{4\eta} \|\nabla u_t\|_2^2 - \omega_2 \frac{\beta - 1}{\beta}\sigma^{\beta/(1-\beta)} \|\nabla u_t\|_{\beta}^{\beta} - \mu \frac{m-1}{m}\delta^{-\frac{m}{m-1}} \|u_t\|_m^m.$$
(3.3)

Next, since u(t) is global and $E(0) \ge d$, then by Theorem 2.3, $E(t) \ge d$, $\forall t \ge 0$. Thus, using the embedding $L^{\alpha}(\Omega) \hookrightarrow L^{2}(\Omega)$ and the inequality

$$z^b \le (z+a)\left(z+\frac{1}{a}\right), \quad z \ge 0, \ 0 < b \le 1, \ a > 0,$$

we obtain

$$\begin{aligned} \|\nabla u\|_{2}^{2} &\leq c \|\nabla u\|_{\alpha}^{\alpha} \\ &= c \left[\|\nabla u\|_{\alpha}^{\alpha} \right]^{2/\alpha} \\ &\leq c \left(1 + \frac{1}{d}\right) \left[\|\nabla u\|_{\alpha}^{\alpha} + d \right] \\ &\leq C \left[\|\nabla u\|_{\alpha}^{\alpha} + E(t) \right], \quad \forall t \geq 0. \end{aligned}$$
(3.4)

By using Lemma 2.1 and (2.2), we get

$$\frac{d}{dt}\left\{F'(t) - \left[\frac{1}{4\eta} + \frac{\beta - 1}{\beta}\sigma^{\frac{-\beta}{\beta - 1}} + \frac{m - 1}{m}\delta^{-\frac{m}{m - 1}}\right]E(t)\right\}$$
$$\geq F''(t) + \frac{\omega_1}{4\eta}\|\nabla u_t\|_2^2 + \omega_2\frac{\beta - 1}{\beta}\sigma^{-\frac{\beta}{\beta - 1}}\|\nabla u_t\|_{\beta}^{\beta} + \mu\frac{m - 1}{m}\delta^{-\frac{m}{m - 1}}\|u_t\|_m^m.$$

Adding and subtracting $p(1-\varepsilon)E(t)$, for $\varepsilon \in (0,1)$, in the right-hand side of the last inequality, and using (3.4) and the Poincaré inequality we obtain

$$\frac{d}{dt} \left\{ F'(t) - \left[\frac{1}{4\eta} + \frac{\beta - 1}{\beta} \sigma^{-\frac{\beta}{\beta - 1}} + \frac{m - 1}{m} \delta^{-\frac{m}{m - 1}} \right] E(t) \right\}$$

$$\geq \|u_t\|_2^2 - \frac{\mu s}{2} \delta^m \|u\|_2^2 - \left[1 + \frac{\omega_2(1 - \theta)}{\alpha} \sigma^\beta \right] \|\nabla u\|_{\alpha}^{\alpha}$$

$$- \left(\omega_1 \eta + \frac{\omega_2 \theta}{2} \sigma^\beta \right) \|\nabla u\|_2^2 + \left[1 - \frac{\mu(1 - s)}{p} \delta^m \right] \|u\|_p^p$$

$$\geq \left[1 + \frac{p}{2}(1 - \varepsilon) \right] \|u_t\|_2^2 - \frac{\mu s}{2} \delta^m \|u\|_2^2 + k(\varepsilon) \|\nabla u\|_{\alpha}^{\alpha}$$

$$- \left(\omega_1 \eta + \frac{\omega_2 \theta}{2} \sigma^\beta \right) \|\nabla u\|_2^2 + \left[\varepsilon - \frac{\mu(1 - s)}{p} \delta^m \right] \|u\|_p^p - p(1 - \varepsilon)E(t)$$

$$\geq \left[1 + \frac{p}{2}(1 - \varepsilon) \right] \|u_t\|_2^2 - \frac{\mu s}{2} \delta^m \|u\|_2^2 + \gamma(\varepsilon) \|\nabla u\|_2^2$$

$$+ \left[\varepsilon - \frac{\mu(1 - s)}{p} \delta^m \right] \|u\|_p^p - [k(\varepsilon) + p(1 - \varepsilon)]E(t)$$

$$\geq \left[1 + \frac{p}{2}(1 - \varepsilon) \right] \|u_t\|_2^2 + \left\{ \gamma(\varepsilon)B - \frac{\mu s}{2} \delta^m \right\} \|u\|_2^2$$

$$+ \left[\varepsilon - \frac{\mu(1 - s)}{p} \delta^m \right] \|u\|_p^p - [k(\varepsilon) + p(1 - \varepsilon)]E(t),$$
(3.5)

where

$$k(\varepsilon) = \frac{1}{\alpha} \left[p(1-\varepsilon) - \alpha - \omega_2(1-\theta)\sigma^\beta \right],$$

$$\gamma(\varepsilon) = \frac{k(\varepsilon)}{C} - \omega_1\eta - \frac{\omega_2\theta}{2}\sigma^\beta,$$

and B is the best constant of Poincaré inequality

$$\|\nabla u\|_2^2 \ge B\|u\|_2^2.$$

Therefore, taking $\eta = \varepsilon$, $\sigma = \varepsilon$,

$$\delta = \left[\frac{p\varepsilon}{\mu(1-s)}\right]^{1/m},$$

setting

$$\gamma_1(\varepsilon) = \frac{1}{4\varepsilon} + \frac{\beta - 1}{\beta} \varepsilon^{-\frac{\beta}{\beta - 1}} + \frac{m - 1}{m} \left(\frac{1 - s}{p\varepsilon}\right)^{-\frac{1}{m - 1}},$$

and substituting in (3.5), we arrive at

$$\frac{d}{dt} \left[F'(t) - \gamma_1(\varepsilon) E(t) \right] \ge \left[1 + \frac{p}{2} (1 - \varepsilon) \right] \|u_t\|_2^2 + \left[\gamma(\varepsilon) B - \frac{ps}{2(1 - s)} \varepsilon \right] \|u\|_2^2 - [k(\varepsilon) + p(1 - \varepsilon)] E(t).$$

By using the Schwarz inequality, we have

$$2\left[1+\frac{p}{2}(1-\varepsilon)\right]^{1/2} \quad \left[\gamma(\varepsilon)B-\frac{ps}{2(1-s)}\varepsilon\right]^{1/2}(u_t,u)$$
$$\leq \left[1+\frac{p}{2}(1-\varepsilon)\right]\|u_t\|_2^2 + \left[\gamma(\varepsilon)B-\frac{ps}{2(1-s)}\varepsilon\right]\|u\|_2^2.$$

Consequently, we obtain

$$\frac{d}{dt} \left[F'(t) - \gamma_1(\varepsilon) E(t) \right] \ge a(\varepsilon)(u_t, u) - \left[k(\varepsilon) + p(1-\varepsilon) \right] E(t) = a(\varepsilon) \left[F'(t) - \gamma_2(\varepsilon) E(t) \right],$$
(3.6)

where

$$a(\varepsilon) = 2 \left[1 + \frac{p}{2} (1 - \varepsilon) \right]^{1/2} \left[\gamma(\varepsilon) B - \frac{ps}{2(1 - s)} \varepsilon \right]^{1/2},$$

$$\gamma_2(\varepsilon) = \frac{k(\varepsilon) + p(1 - \varepsilon)}{a(\varepsilon)}.$$

Since

$$\gamma(\varepsilon)B - \frac{ps}{2(1-s)}\varepsilon \to \begin{cases} \frac{B(p-\alpha)}{\alpha C} > 0 & \text{as } \varepsilon \to 0^+ \\ -\left[\frac{\alpha + \omega_2(1-\theta)}{\alpha C} + \omega_1 + \frac{\omega_2\theta}{2}\right]B - \frac{ps}{2(1-s)} < 0 & \text{as } \varepsilon \to 1^-, \end{cases}$$

then, there exists $\varepsilon_* \in (0, 1)$, such that

$$a(\varepsilon_*) = 0$$
 and $a(\varepsilon) > 0$, $\forall \varepsilon \in (0, \varepsilon_*)$.

Hence, we have

$$\gamma_1(\varepsilon) - \gamma_2(\varepsilon) \to \begin{cases} +\infty & \text{as } \varepsilon \to 0^+ \\ -\infty & \text{as } \varepsilon \to \varepsilon_*^-. \end{cases}$$

Therefore, there exists $\varepsilon_0 \in (0, \varepsilon_*)$, such that $\gamma_1(\varepsilon_0) = \gamma_2(\varepsilon_0) > 0$. So, by setting

$$L(t) = F'(t) - \gamma_1(\varepsilon_0)E(t),$$

$$M = \gamma_1(\varepsilon_0),$$
(3.7)

and by using (2.3), we obtain

$$L(0) = (u_t(0), u(0)) - \gamma_1(\varepsilon_0)E(0)$$

> $(u_t(0), u(0)) - ME(0) > 0.$

Moreover, with this choice of ε_0 , (3.6) becomes

$$\frac{d}{dt}L(t) \ge a(\varepsilon_0)L(t),$$

which gives

$$L(t) \ge L(0)e^{a(\varepsilon_0)t}, \quad \forall t \ge 0,$$

and hence

$$F'(t) \ge L(0)e^{a(\varepsilon_0)t}, \quad \forall t \ge 0.$$

By integrating this last inequality over (0, t), we get

$$\|u(t)\|_{2}^{2} = 2F(t) \ge 2F(0) + 2\frac{L(0)}{a(\varepsilon_{0})} \left[e^{a(\varepsilon_{0})t} - 1\right], \quad \forall t \ge 0.$$
(3.8)

On the other hand, by using Hölder's inequality and (2.2), we have

$$\begin{aligned} \|u(t)\|_{2} &\leq \|u(0)\|_{2} + \int_{0}^{t} \|u_{\tau}(\tau)\|_{2} d\tau \\ &\leq \|u(0)\|_{2} + C \int_{0}^{t} \|u_{\tau}(\tau)\|_{m} d\tau \\ &\leq \|u(0)\|_{2} + Ct^{\frac{m-1}{m}} \int_{0}^{t} \|u_{\tau}(\tau)\|_{m}^{m} d\tau \\ &\leq \|u(0)\|_{2} + Ct^{\frac{m-1}{m}} \int_{0}^{t} \frac{-1}{\mu} \frac{dE(\tau)}{d\tau} d\tau \\ &\leq \|u(0)\|_{2} + Ct^{\frac{m-1}{m}} \left[\frac{E(0) - E(t)}{\mu}\right]^{1/m} \\ &\leq \|u(0)\|_{2} + C \left[\frac{E(0)}{\mu}\right]^{1/m} t^{\frac{m-1}{m}}, \end{aligned}$$

which clearly contradicts (3.8).

4. Lower bound for the blow-up time

In this section, we give a lower bound for the blow-up time $T_{\rm max}.$ To this end, we define

$$G(t) := \frac{1}{p} \| u(t) \|_p^p.$$

Theorem 4.1. Let u be the solution of (1.1), and assume that

$$\left\{ \begin{array}{ll} 2 \alpha, \\ 2$$

Then

$$T_{\max} \ge \int_{G(0)}^{+\infty} \left\{ \tau + A_1 \tau^{\frac{2}{\alpha}(p-1)} + A_2 \right\}^{-1} d\tau,$$

where A_1 and A_2 are positive constants to be determined later in the proof.

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Proof. By using inequality (2.1), we have

$$\frac{1}{2} \|u_t\|_2^2 + \frac{1}{\alpha} \|\nabla u\|_{\alpha}^{\alpha} = E(t) + \frac{1}{p} \|u(t)\|_p^p \le E(0) + G(t).$$
(4.1)

Next, using the Schwarz inequality, the Sobolev-type inequality

$$\|u\|_q \le C_q \|\nabla u\|_{\alpha}, \quad \forall q \in [1, \alpha^*], \quad \forall u \in W_0^{1, \alpha}(\Omega),$$

$$(4.2)$$

inequality (4.1) yields

$$G'(t) = (|u|^{p-2}u, u_t)$$

$$\leq \frac{1}{2} ||u_t||_2^2 + \frac{1}{2} ||u||_{2(p-1)}^{2(p-1)}$$

$$\leq \frac{1}{2} ||u_t||_2^2 + \frac{C_{2(p-1)}^{2(p-1)}}{2} ||\nabla u||_{\alpha}^{2(p-1)}$$

$$\leq E(0) + G(t) + \frac{C_{2(p-1)}^{2(p-1)}}{2} [\alpha E(0) + \alpha G(t)]^{\frac{2}{\alpha}(p-1)}.$$
(4.3)

From (4.3) and Jensen's inequality, we obtain the differential inequality

$$G'(t) \le G(t) + A_1 \left[G(t) \right]^{\frac{2}{\alpha}(p-1)} + A_2, \tag{4.4}$$

with

$$A_1 = C_*^{2(p-1)} 2^{\frac{2}{\alpha}(p-1)-2} \alpha^{\frac{2}{\alpha}(p-1)} \quad \text{and} \quad A_2 = E(0) + A_1 \left[E(0) \right]^{\frac{2}{\alpha}(p-1)}.$$

Hence, we get

$$T_{\max} \ge \int_0^{T_{\max}} \left\{ G(s) + A_1 \left[G(s) \right]^{\frac{2}{\alpha}(p-1)} + A_2 \right\}^{-1} G'(s) ds.$$

Since $\lim_{t \to T_{\max}^-} G(t) = +\infty$, so the previous inequality implies

$$T_{\max} \ge \int_{G(0)}^{+\infty} \left\{ \tau + A_1 \tau^{\frac{2}{\alpha}(p-1)} + A_2 \right\}^{-1} d\tau.$$

In the next theorem, when $n > \alpha$, the upper bound for p is enlarged. We define

$$H(t) := \frac{1}{\sigma} \|u(t)\|_{\sigma}^{\sigma},$$

where $\sigma = \frac{\alpha(n-2)+2n}{2(n-\alpha)}$. Then, we have

Theorem 4.2. Let u be the solution of (1.1), and assume that

$$\frac{\alpha(n-2)+2n}{2(n-\alpha)}
(4.5)$$

Then

$$T_{\max} \ge \int_{H(0)}^{+\infty} \left\{ B_1 \tau^{b_1} + B_2 \tau^{b_2} + B_3 \right\}^{-1} d\tau,$$

where B_1, B_2, B_3, b_1 and b_2 are positive constants to be determined later in the proof.

Proof. By using inequality (2.1), we have

$$\frac{1}{2} \|u_t\|_2^2 + \frac{1}{\alpha} \|\nabla u\|_{\alpha}^{\alpha} = E(t) + \frac{1}{p} \|u(t)\|_p^p \le E(0) + \frac{1}{p} \|u(t)\|_p^p.$$
(4.6)

Using the Schwarz inequality, the Sobolev-type inequality (4.2), with $q = \alpha^*$, and inequality (4.6) we get

$$H'(t) = (|u|^{\sigma-2}u, u_t)$$

$$\leq \frac{1}{2} ||u_t||_2^2 + \frac{1}{2} ||u||_{2(\sigma-1)}^{2(\sigma-1)}$$

$$\leq \frac{1}{2} ||u_t||_2^2 + \frac{C_*^{\alpha^*}}{2} ||\nabla u||_{\alpha}^{\alpha^*}$$

$$\leq E(0) + \frac{1}{p} ||u||_p^p + \frac{C_*^{\alpha^*}}{2} \left[\alpha E(0) + \frac{\alpha}{p} ||u||_p^p \right]^{\frac{n}{n-\alpha}}.$$
(4.7)

Next, the interpolation inequality, the Sobolev inequality and Young's inequality give

$$\begin{aligned} \|u\|_{p}^{p} &\leq \|u\|_{\alpha^{*}}^{\theta_{p}} \cdot \|u\|_{\sigma}^{(1-\theta)p}, \qquad \theta = \frac{\alpha^{*}(p-\sigma)}{p(\alpha^{*}-\sigma)}, \\ &\leq C_{*}^{\theta_{p}} \|\nabla u\|_{\alpha}^{\theta_{p}} \cdot \|u\|_{\sigma}^{(1-\theta)p}, \\ &\leq \frac{1}{\alpha} \|\nabla u\|_{\alpha}^{\alpha} + B\|u\|_{\sigma}^{r}, \end{aligned}$$

$$(4.8)$$

where

$$B = C_* \left(1 - \frac{\theta p}{\alpha} \right) (p\theta C_*)^{\frac{p\theta}{\alpha - p\theta}} \quad \text{and} \quad r = \frac{\alpha p(1 - \theta)}{\alpha - \theta p}.$$

Note that in virtue of (4.5), we have $\alpha > \theta p$. Hence, by (2.1) we have

$$\|u\|_{p}^{p} \le E(0) + \frac{1}{p} \|u\|_{p}^{p} + B\|u\|_{\sigma}^{r},$$
(4.9)

which gives

$$\frac{1}{p} \|u\|_p^p \le \frac{1}{p-1} \left(E(0) + B \|u\|_{\sigma}^r \right).$$

Inserting this last inequality in (4.7), and using Jensen's inequality, we obtain

$$H'(t) \leq \frac{pE(0)}{p-1} + \frac{B}{p-1} \|u\|_{\sigma}^{r} + \frac{C_{*}^{\alpha^{*}}}{2} \left[\frac{\alpha pE(0)}{p-1} + \frac{\alpha B}{p-1} \|u\|_{\sigma}^{r} \right]^{\frac{n}{n-\alpha}} = B_{1} \left(H(t)\right)^{b_{1}} + B_{2} \left(H(t)\right)^{b_{2}} + B_{3},$$
(4.10)

where

$$B_1 = \frac{B\sigma^r}{p-1}, \quad B_2 = \frac{C_*^{\alpha^*}}{2} 2^{\frac{\alpha}{n-\alpha}} \left[\frac{\alpha B\sigma^r}{p-1} \right]^{\frac{n}{n-\alpha}},$$
$$B_3 = \frac{pE(0)}{p-1} + \frac{C_*^{\alpha^*}}{2} 2^{\frac{\alpha}{n-\alpha}} \left[\frac{\alpha pE(0)}{p-1} \right]^{\frac{n}{n-\alpha}},$$
$$b_1 = \frac{r}{\sigma}, \quad b_2 = \frac{rn}{\sigma(n-\alpha)}.$$

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Finally, integrating inequality (4.10) over $(0, T_{\text{max}})$ we get

$$T_{\max} \ge \int_0^{T_{\max}} \left\{ B_1 \left(H(s) \right)^{b_1} + B_2 \left(H(s) \right)^{b_2} + B_3 \right\}^{-1} H'(s) ds,$$

and so

$$T_{\max} \ge \int_{H(0)}^{+\infty} \left\{ B_1 \tau^{b_1} + B_2 \tau^{b_2} + B_3 \right\}^{-1} d\tau.$$

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