

THE "ZIG-ZAG" CYLINDER RULE

MIRCEA V. DIUDEA^a

^a Faculty of Chemistry and Chemical Engineering
Babes-Bolyai University, 3400 Cluj, Romania

ABSTRACT. Spectral data revealed some cluster properties of tubulenes. The "zig-zag" cylinder rule on the π -electronic structure states that properly closed shell exists in two complementary series of fullerenes derived from the $13k$ cluster, $k = 4m; 4m+2$, $m = 1, 2, \dots$, by adding a $3kn$, $n = 1, 2, \dots$, nanotube distancer between the two zig-zag ended caps. Semiempirical calculations support this finding.

INTRODUCTION

A fullerene is, according to a classical definition, an all-carbon molecule consisting entirely of pentagons (exactly 12) and hexagons ($N/2-10$).¹ Besides the well known near-spherical fullerenes, open ended nanotubes, capped tubules and tori have aroused both theoretical and experimental interest.²⁻¹⁰ Non-classical fullerene extensions to include rings of other sizes have been considered.^{11,12} Multi elemental large cages have also been studied.¹³

In simple π -only Hückel theory, the energy of the i^{th} molecular orbital $E_i = \alpha + \lambda_i \beta$ is calculated on the ground of the adjacency matrix associated to the molecular hydrogen depleted graph. Systematic studies on this matrix and their eigenvalue spectra (*i.e.*, the decreasing sequence of eigenvalues λ_i), provided some magic number rules for the stability of molecules, such as the oldest Hückel¹⁴ $4n+2$ rule for aromatic rings and the more recent $60 + 6m$ ($m \neq 1$) *leapfrog* rule¹⁵ for the properly closed fullerenes (see also ref. 16).

The π -electronic shells of neutral fullerenes are classified, function of their eigenvalue spectra, as:¹⁷ (i) *properly closed*, PC, when $\lambda_{N/2} > 0 \geq \lambda_{N/2+1}$; (ii) *pseudo-closed*, PSC, in case $\lambda_{N/2} > \lambda_{N/2+1} > 0$; (iii) *meta-closed*, MC, with $0 \geq \lambda_{N/2} > \lambda_{N/2+1}$ and (iv) *open*, OP, when the $N/2^{\text{th}}$ (HOMO) and $N/2+1^{\text{th}}$ (LUMO) molecular orbitals are degenerate, $\lambda_{N/2} = \lambda_{N/2+1}$. The bandgap is taken as the absolute value of the difference $E_{\text{HOMO}} - E_{\text{LUMO}}$. The most frequent case is that of the pseudo-closed shell, since the number of positive eigenvalues is, in general, larger than that of the negative ones, $n_+ \geq n_-$.¹⁸

The other rule predicting closed-shell fullerenes is the *cylinder rule*.^{19,20} Fullerenes with k -fold cylindrical symmetry, of general formula $C_{N,k-V[2k,n]-[6]}$, have a closed shell at each nuclearity $N = 2k(7 + 3m)$, $m = 0, 1, 2, \dots$, ($k = 4$ to 7 in this paper). These cages have a non-degenerate non-bonding orbital (NBO) LUMO separated by a gap from HOMO; in this case $n_+ = 1 + n_-$. Exceptions exist, e.g., the first term of series $k = 6$, $C_{72,6-V[2k,1]-[6]}$, has LUMO triply degenerate, and the first term of series $k = 7$, $C_{84,7-V[2k,1]-[6]}$, has its NBO not the LUMO (actually is $\lambda_{N/2+3}$).

The cylinder rule, better called the "armchair" cylinder rule, can equivalently be written, in distancing tube dimension, as: $n = 1 + 3m$.

Recall that, our notations for nanotubes $\text{TUVC}_6[c,n]$ and $\text{TUHC}_6[c,n]$ correspond to the "armchair" ($c/2, c/2$) and "zig-zag" ($c/2, 0$), respectively.

This paper describes a further investigation of fullerenes with the aim of finding some novel rules of their electronic stability.

TUBULENES BY "ZIG-ZAG" NANOTUBE DISTANCER

A spherical fullerene, say C_{60} (I_h), can provide a cap suitable for joining an armchair nanotube (i.e., a TUVC_6); the cap, having a k -fold polar ring, will be denoted by $C_{N,k-V[2k,0]}$ with N being the number of atoms in the parent fullerene. In a fullerene cluster the polar ring can vary, in our work most often $k = 4, 5, 6, 7$ (with the main term at $k = 5$). The cap $C_{N,k-V[2k,0]}$ can be used in generating a cluster of tubulenes by using a TUV distancer. The procedure is illustrated in Figure 1. This is the cluster $C_{N,k-V[2k,n]-[6]}$, obeying the armchair cylinder rule (initially defined for $k = 5; 6$).¹⁹

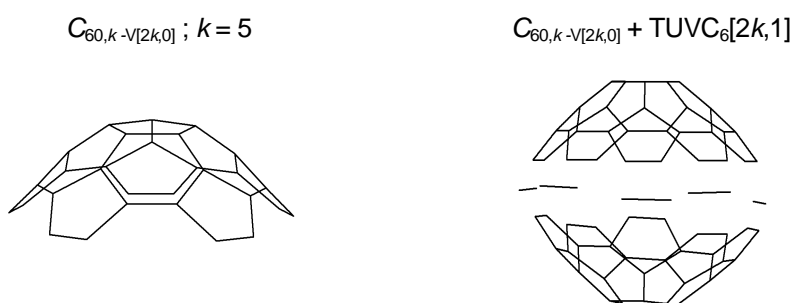
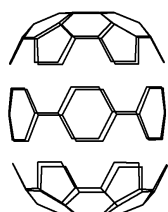


Figure 1. "Armchair" tubulene construction

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$$C_{72,k-V[2k,0]} + \text{TUVC}_6[2k,3]; k = 6$$



$$C_{18k}$$

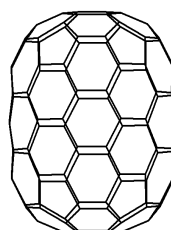
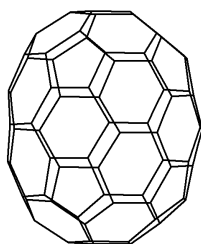


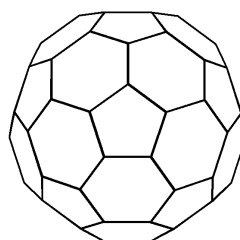
Figure 1. (continued)

The first term of the sub-cluster $C_{60,5-V[10,1]-[6]}$ is just the isolated fullerene C_{70} (see below). Its topology is given as the spiral code.^{21,22} We amended this code by specifications for the mode of inserting the new polygon, either by an edge (e) or by an angle (a).

$$C_{60,5-V[10,1]-[6]} = C_{70}$$



$$C_{70}; (\text{top})$$



$$k, 6k, (a5, e6)k, (e6, e6)k, (a5, e6)k, k$$

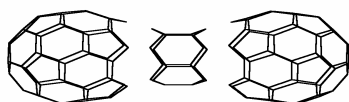
$$k = 4; 5; 6; 7$$

$$\text{Cluster } C_{14k}, k = 4; 5; 6; 7$$

$$N = 14k, 56; 70; 84; 98$$

A different cap, denoted $C_{N,k-H[2k,0]}$, is that favoring the coupling with a zig-zag nanotube (*i.e.*, TUHC_6).²⁰ Figure 2 gives such an example.

$$2(C_{60,5-H[10,0]}) + \text{TUHC}_6[10,2]$$



$$C_{60,5-H[10,3]-[7,6,7]}$$

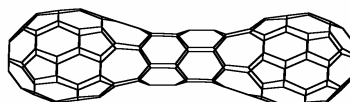
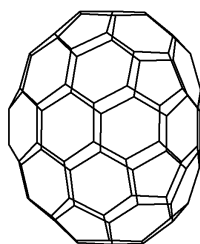


Figure 2. "Zig-zag" tubulene construction

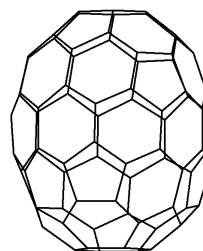
A tubulene $C_{N,k-V[2k,6]-[6]}$ can be viewed as originating in the coalescence product of two spherical fullerenes, namely the $C_{N,k-H[2k,1]-[7]}$, from which it results by a series of Stone-Wales^{23,24} edge flipping. Similar coalescence reactions between spherical cages and nanotubes have been considered.²⁵

Resuming to the nomenclature, the *C* letter (*i.e.*, cage) is followed by the number of atoms *N* of the parent fullerene, having a *k*-fold polar ring, the cap being attached to a *V/H*[*c,n*] tube (with a cross-section of *c* atoms, and *n* atom rows distancing the two caps; *n* = 0 for the cap, only). The numbers in the last brackets denote the tiling polygons in the region between the two caps.

A quite different cap $C_{N,k-H[c,0]}$ is provided by the isolated fullerene C_{78} , the two of its IP (*i.e.*, isolated pentagon) isomers being given below.

 $C_{78}a$; C_S 

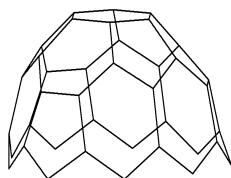
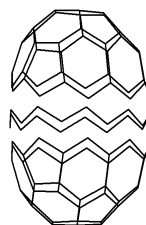
$k, (5, 6)k/2, (a6, a6, e5)k/2, (a6)k,$
 $(e5, a6, a6)k/2, (6, 5)k/2, k$
 $k = 6$

 $C_{78}c$; C_1 

Cluster C_{13k} , $k = 4; 6; 8; 10$

$N = 13k; 52; 78; 104; 130$

The formal building of the related tubulenes is illustrated in Figure 3.

 $C_{78,k-H[3k,0]}$; $k = 6$  $C_{78,k-H[3k,0]} + \text{TUHC}_6[3k,1]$  $C_{78,k-H[3k,0]} + \text{TUHC}_6[3k,2]$ $C_{78,k-H[3k,2]-[6]} = C_{19k}$; $k = 6$

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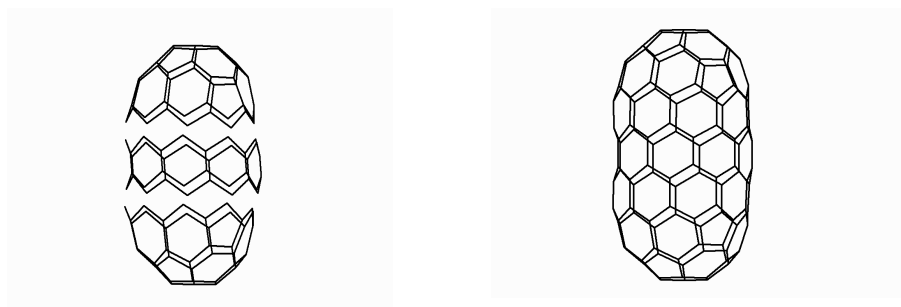


Figure 3. "Zig-zag" tubulenes derived from C_{78}

SPECTRAL PROPERTIES OF THE NOVEL CLUSTER

The zig-zag tubulenes derived from the cluster C_{13k} form a properly closed shell tubulene cluster with the general formula $C_{13k,k-H[3k,n]-[6]}$ and the following construction $N = 13k + 3kn$; $k = 4m + 2z$; $m = 1, 2, \dots$; $z = 0, 1$; $n = 1, 2, \dots$. It shows the same bandgap for a given z value at each positive n . The last parameter is just the n -dimension of the distancing tube and it discriminates between the twin odd and even sub-clusters (see Figures 4 and 5).

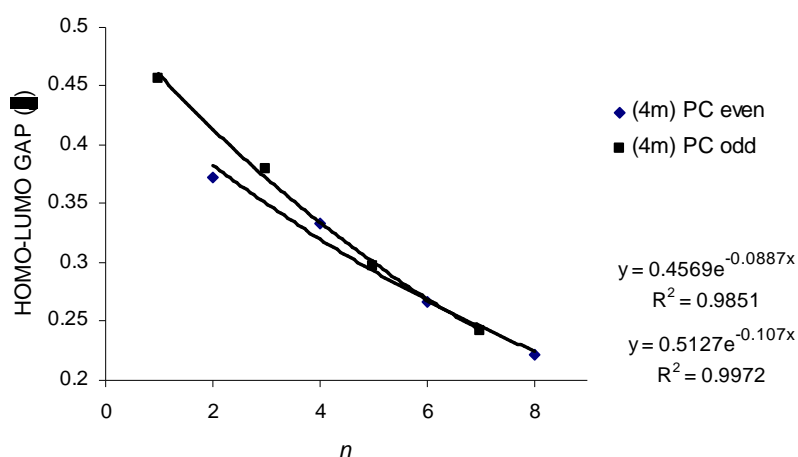


Figure 4. The plot of the spectral gap vs. n , in the series $k = 4m$

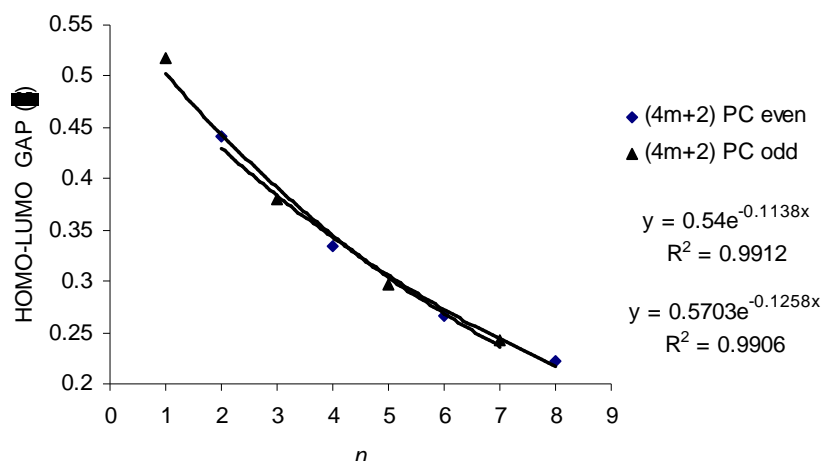


Figure 5. The plot of the spectral gap vs. n , in the series $k = 4m+2$

In the opposite to the 1st (armchair) cylinder rule, the terms for $n = 0$ do not (properly) belong to the new cluster. In the series $k = 4m$, LUMO is an NBO (not encountered at higher n -values) while in the series $k = 4m+2$, the gaps are different for different m -values. The identity of the gaps appears only at $n > 0$. Figure 6 gives the repeating units of the twin sub-clusters, in the so-called geodesic projection.

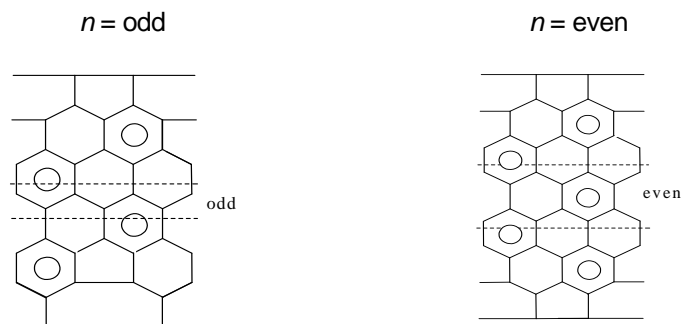


Figure 6. Geodesic projection of the repeating units of the twin sub-clusters with the general formula $C_{13k,k-H[3k,n]-[6]}$

The novel cluster $C_{13k,k-H[3k,n]-[6]}$, of properly closed shell fullerenes, having $N = 13k + 3kn$; $n = 1, 2, \dots$ and equal number of positive and negative eigenvalues $n_+ = n_-$, we call *the zig-zag cylinder rule*, to specify the type of the tube distancing the two caps. This cluster does not superimpose over the

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leapfrog^{15,18,19} (properly closed-shell) fullerenes with the exception of the cages having $k = 6$. The spectral data are supported by semiempirical calculations (see Table 1).

In going from one sub-cluster to the other, complementary (twisted) series of pseudo-closed shell tubulenes appear. Figure 7 shows the plots of PSC and PC series corresponding to $k = 6$ (viewed as single odd&even series). Clearly, the gap is deeper for the PC series, suggesting a higher kinetic stability of the corresponding tubulenes. Note that the term $n = 0$ in PSC is just the $C_{78}C_1$ (see above).

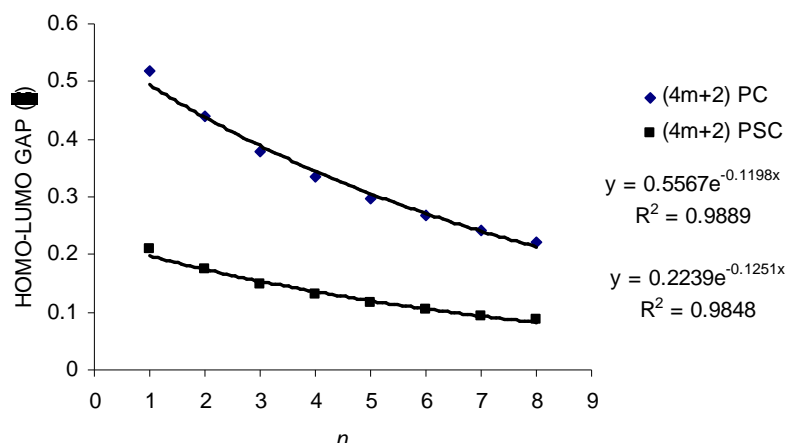


Figure 7. Comparison between the PSC and PC series corresponding to $k = 6$

In the case of objects with formula $C_{N,k-H[2k,n]-[7,6,7]}$ (see Figure 2) the open shell appears, starting with the first row of hexes ($n = 2$) distancing the two caps. The two frontier orbitals degenerate within the positive domain of eigenvalue for $k = 4$, while for higher k they are located either in the positive or negative domain (as in some tetrahedral fullerenes, with $N \geq 628$).¹⁸ The most important exception is $C_{72,6-H[12,2]-[7,6,7]}$ ($N = 156$), with a meta-closed shell. The "accidental" gap is, however, very thin (-0.026439; -0.028345). This is the only case of MC shell reported in spherical fullerenes (more exactly, fullerenes of genus zero).¹⁸

SEMIEMPIRICAL CALCULATIONS

The semiempirical calculations were calculated with the PM3 Hamiltonian (by HyperChem software package).²⁶ Data are given in Table 1.

Fowler²⁷ has found that geometric instability appears in fullerenes at a "small" gap of ca. 5 eV, in semiempirical calculation, or 0.4 $|\beta|$ in simple Hückel calculation. At least the first terms, and particularly in the series $k = 4m+2$

(Table 1, entries 4-6 and 10-12), show favorable bandgap and relatively low heat of formation (per atom). Compare these data with those for C_{60} : HF = 13.512; PM3 Gap = 6.594 and find that we are in the same domain. Besides the isolated C_{78} (entry 4), some other tubulenes can be candidates to the status of real molecules. A cage having a 10-fold face could appear strange but it just obeys the Hückel $4m+2$ aromatic rule.

Table 1.

Data for the tubulenes $C_{13k,k-H[3k,n]-[6]}$

	Cage k, n	N	Sym.	PM3	PM3	Spectral Data			Shell
				HF/at. (kcal/mol)	GAP (eV)	$\lambda_N / 2$	$\lambda_N / 2+1$	GAP ($ \beta $ units)	
1	4; 0	52	D_{2h}	21.585	5.533	0.2564	0	0.2564	PC
2	4; 1	64	S_4	19.271	5.783	0.3789	-0.0774	0.4563	PC
3	4; 2	76	D_{2h}	18.162	5.469	0.3260	-0.0470	0.3731	PC
4	6; 0	78	C_s	12.294	6.083	0.5157	-0.1176	0.6333	PC
5	6; 1	96	D_{3d}	11.015	5.564	0.4329	-0.0844	0.5173	PC
6	6; 2	114	C_s	10.381	5.195	0.3688	-0.0721	0.4409	PC
7	8; 0	104	C_{4h}	11.453	5.730	0.2564	0	0.2564	PC
8	8; 1	128	D_{4d}	10.158	5.497	0.3789	-0.0774	0.4563	PC
9	8; 2	152	C_{4h}	9.390	5.112	0.3260	-0.0470	0.3731	PC
10	10; 0	130	C_{5h}	12.518	5.998	0.4579	-0.0874	0.5453	PC
11	10; 1	160	D_{5d}	11.875	5.512	0.4329	-0.0844	0.5173	PC
12	10; 2	190	C_{5h}	10.715	5.115	0.3688	-0.0721	0.4409	PC

CONCLUSIONS

The zig-zag cylinder rule, first reported in this paper, presents a cluster with the general formula $C_{13k,k-H[3k,n]-[6]}$, of properly closed shell tubulenes, having $N = 13k + 3kn$; $n = 1, 2, \dots$. The number n is just the number of atom rows in the tube distancing the two caps. It shows the same bandgap for either odd or even k value at each positive n . The semiempirical calculations support the idea of relatively stable molecules, possible appearing in the soot of the vaporized graphite.

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